## **Optical Metrology**

Lecture 6: Moiré and Triangulation

# Sinusoidal gratings

Often physical gratings are square-wave function, but we can describe them by sinusoidal gratings

$$t_1(x, y) = a + a \cos\left(\frac{2\pi}{p}x\right)$$

where p is the grating period and where  $0 < a < \frac{1}{2}$ .

These gratings become modulated

$$t_2(x, y) = a + a\cos 2\pi \left(\frac{x}{p} + \psi(x)\right)$$
(7.2)

 $\psi(x)$  is the modulation function and is equal to the displacement of the grating lines from its original position divided by the grating period

$$\psi(x) = \frac{u(x)}{p} \tag{7.3}$$

where u(x) is the displacement.

# Sinusoidal gratings

When the two gratings given by Equations (7.1) and (7.2) are laid in contact, the resulting transmittance *t* becomes the product of the individual transmittances, viz.

$$t(x, y) = t_{1}t_{2}$$

$$= a^{2} \left\{ 1 + \cos \frac{2\pi}{p} x + \cos 2\pi \left[ \frac{x}{p} + \psi(x) \right] + \frac{1}{2} \cos 2\pi \left[ \frac{2x}{p} + \psi(x) \right] + \frac{1}{2} \cos 2\pi \psi(x) \right\}$$
(7.4)

- The first three terms represent the original gratings, the fourth term the second grating with doubled frequency,
- The fifth term depends on the modulation function only. It is this term which describes the moiré pattern.

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(7.4)

a maximum resulting in a bright fringe whenever

$$\psi(x) = n$$
, for  $n = 0, \pm 1, \pm 2, \pm 3, \dots$  (7.6)

and minima (dark fringes) whenever

$$\psi(x) = n + \frac{1}{2}, \text{ for } n = 0, \pm 1, \pm 2, \pm 3, \dots$$
 (7.7)

Both grating  $t_1$  and  $t_2$  could be phase-modulated by modulation functions  $\psi_1$  and  $\psi_2$  respectively. Then  $\psi(x)$  in Equations (7.6) and (7.7) has to be replaced by

$$\psi(x) = \psi_1(x) - \psi_2(x) \tag{7.8}$$

# Moiré

- The figure is an illustration of two interfering plane waves.
- Two gratings that lie in contact, with a small angle between the grating lines.
- As a result, we see a fringe pattern of much lower frequency than the individual gratings.
- This is an example of the moiré effect and the resulting fringes are a moiré pattern.



## Moiré

# Transmittance is equal to the product $t(x, y) = t_1 t_2$

#### Interfringe distance

 $d = \frac{p}{2\sin\frac{\alpha}{2}}$ 

This can be applied for measuring  $\alpha$  by measurement of d.





When measuring in-plane deformations a grating is attached to the test surface. When the surface is deformed, the grating will follow the deformation and will therefore be given by Equation (7.2). The deformation u(x) will be given directly from Equation (7.3):

$$u(x) = p\psi(x) \tag{7.10}$$

To obtain the moiré pattern, one may apply one of several methods (Post 1982; Sciammarella 1972, 1982):

(1) Place the reference grating with transmittance  $t_1$  in contact with the model grating with transmittance  $t_2$ . The resulting intensity distribution then becomes proportional to the product  $t_1 \cdot t_2$ .

To obtain the moiré pattern, one may apply one of several methods (Post 1982; Sciammarella 1972, 1982):

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We obtain:

 $I(x) = I_0 + I_1 \cos 2\pi \psi(x) + \text{terms of higher frequencies}$ 

A DC term *I*<sup>0</sup> and a term containing the modulation function. When applying low frequency gratings -> OK. High frequency -> Low contrast Moiré fringes.

We can filter optically or digitally to obtain:

maximum whenever  $\psi(x) = n$ , for n = 0, 1, 2, ...

minimum whenever  $\psi(x) = n + \frac{1}{2}$ , for n = 0, 1, 2, ...

According to Equation (7.10) this corresponds to a displacement equal to

$$u(x) = np$$
 for maxima (7.13a)

$$u(x) = (n + \frac{1}{2})p$$
 for minima (7.13b)

Intensity distribution of Moiré pattern, displacement and strain.



By orienting the model grating and the reference grating along the y-axis, we can in the same manner find the modulation function  $\psi_y(y)$  and the displacement v(y) in the y-direction.  $\psi_x(x)$  and  $\psi_y(y)$  can be detected simultaneously by applying crossed gratings, i.e. gratings of orthogonal lines in the x- and y-directions. Thus we also are able to calculate the strains

$$\varepsilon_x = p \frac{\partial \psi_x}{\partial x} \tag{7.14a}$$

$$\varepsilon_y = p \frac{\partial \psi_y}{\partial y}$$
 (7.14b)

$$\gamma_{xy} = p \left[ \frac{\partial \psi_x}{\partial y} + \frac{\partial \psi_y}{\partial x} \right]$$
(7.14c)

### Measurement of Out-of-Plane Deformations. Contouring

An effect where moiré fringes are formed between a grating and its own shadow.

P0 is projected onto P1, then by viewing is projected to P2 on the grating.

Equivalent to displacement:

$$u = u_1 + u_2 = h(x, y)(\tan \theta_1 + \tan \theta_2)$$

h(x, y) is the height difference



Shadow Moiré

Figure 7.3 Shadow moiré

#### Measurement of Out-of-Plane Deformations. Contouring

This corresponds to a modulation function

$$\psi(x) = \frac{u}{p} = \frac{h(x, y)}{p} (\tan \theta_1 + \tan \theta_2)$$

A bright fringe is obtained when

$$\psi(x) = n$$
, for  $n = 0, 1, 2, ...,$   
 $h(x, y) = \frac{np}{\tan \theta_1 + \tan \theta_2}$ 



Shadow Moiré

Figure 7.3 Shadow moiré

**Dark fringe** 
$$h(x, y) = \frac{(n + \frac{1}{2})p}{\tan \theta_1 + \tan \theta_2}$$

# Triangulation

beam is incident on a diffusely scattering surface under an angle  $\theta_1$ . The resulting light spot on the surface is imaged by a lens onto a detector D. The optical axis of the lens makes an angle  $\theta_2$  to the surface normal. Assume that the object moves a distance *s* normal to its surface. From the figure, using simple trigonometric relations, we find that



Figure 7.13 Triangulation probe